



Thermoacoustic elastography: recovery of bulk elastic modulus of heterogeneous media using tomographically measured thermoacoustic measurements

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Background: Tissue mechanical parameters such as elasticity are of great significance for the assessment of biological histopathological and physiological characteristics. Here, we propose a new approach called thermoacoustic elastography (TAE) for imaging tissue elastic modulus.

Methods: Central to TAE is an image reconstruction algorithm that allows the recovery of both microwave energy loss and elastic modulus distributions. The algorithm is first evaluated using simulated data under various practical scenarios, including a varied range of microwave energy loss and elastic modulus between the heterogeneity and background region, different noise levels, and multiple targets. The feasibility of the proposed TAE was then validated by imaging the elastic modulus distribution of agar phantoms with various elastic modulus and microwave energy loss.

Results: The results from both the simulated and phantom experiments show that the recovered elastic modulus by TAE agree well with the exact values, having an average error of less than 12.74%.

Conclusions: The findings from this study suggest that TAE provides a new addition to the family of elasticity imaging and may have broad application prospects, such as cirrhosis and atherosclerosis detection.

Keywords: Bulk elastic modulus; thermoacoustic imaging; medical imaging

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Introduction

Elastic modulus of tissue (1,2) is an important parameter in disease diagnosis, which can characterize tissue as soft or hard. Elastography for measuring tissue elastic modulus (2-9) is widely used as a supplementary diagnostic method for early detection and diagnosis of various diseases. At present, elastography has been implemented based on the contrast mechanism of ultrasound (US) imaging (3,10,11), optical coherence tomography (OCT) (4,12), magnetic resonance imaging (5,13) (MRI), or photoacoustic imaging (PAI) (6-9,14,15). These elasticity imaging techniques have been applied for the detection of liver fibrosis (16), breast

cancer (17), and prostate cancer (18).

Here, we add a new method to the family of elastography based on the contrast mechanism of thermoacoustic tomography (TAT). TAT (19-25) is an emerging imaging technique based on the thermoacoustic (TA) effect, and provides high microwave contrast, high spatial resolution, and large penetration depth. Thermoacoustic elastography (TAE) has a wide range of potential medical applications, e.g., in the detection of cirrhosis, atherosclerosis, and cancers (25-27). In TAE, we reconstruct the bulk elastic modulus distribution by solving the TA wave equation based on the finite element method (FEM) without

internal or external static stress or shear wave. We calculate acoustic pressure through the initial microwave energy loss and bulk elastic modulus distributions. The microwave energy loss and elastic modulus distributions are then iteratively updated until the error between the calculated and measured acoustic pressures meets a pre-set tolerance, or the number of iterations reaches a pre-set upper limit. Compared with optical coherence elastography and photoacoustic elastography, TAE has superior imaging depth. Compared to magnetic resonance elastography, TAE has higher resolution. Our TAE approach is tested and validated using both simulated and phantom experiments.

Methods

Image reconstruction algorithm

Key to the realization of TAE is an image reconstruction algorithm based on the FEM that simultaneously reconstructs tissue microwave energy loss and elastic modulus from tomographically measured TA signals.

Using the Newton's law of motion, equation of continuity, and thermal elastic equation, the following time-domain TA wave equation can be derived (28):

$$\nabla^2 p(r,t) - \frac{\rho}{k} \frac{\partial^2}{\partial t^2} p(r,t) = -\frac{\beta}{c_p} \frac{\partial}{\partial t} H(r,t) \quad [1]$$

where p is the pressure wave, ρ is the mass density, t is the time, r is the spatial variable, K is the bulk elastic modulus, C_p is the specific heat capacity at constant pressure, and β is the thermal coefficient of volume expansion. H is the microwave excitation source term, which can be written as $H = \Psi(r)I(t)$, where $I(t)$ is the temporal illumination function and $\Psi(r)$ is the absorbed energy density. To obtain Eq. [1], a homogeneous elastic reference medium is assumed with density, $\rho = \rho_0$, and the bulk elastic modulus is denoted as $K = \rho v^2$, where v is the acoustic velocity. If we let $O = \frac{\rho_0}{K}$, and expand p and O as the sum of coefficients multiplied by a set of the basis function ψ_j and ψ_k —i.e., $p = \sum \psi_j p_j$ and $O = \sum \psi_k O_k$ —then the finite-element discretization of Eq. [1] is expressed as (29)

$$\sum_{j=1}^N p_j \left[\int_s \nabla \psi_i \cdot \nabla \psi_j ds \right] + \sum_{j=1}^N \ddot{p}_j \left[\int_s \psi_i \psi_j \psi_k O_k ds \right] - \oint_l \psi_i \nabla p \cdot \hat{n} dl = \int_s \frac{\beta \Psi}{C_p} \frac{\partial I}{\partial t} \psi_i ds \quad [2]$$

where \hat{n} is the unit normal vector.

The Bayliss-Turkel radiation boundary conditions are employed here (28):

$$\nabla p \cdot \hat{n} = -\frac{1}{v_0} \frac{\partial p}{\partial t} - \frac{p}{2r} \quad [3]$$

In both the forward and the inverse calculations, the unknown coefficients Ψ need to be expanded in a similar fashion to p , as a sum of unknown parameters multiplied by a known spatially varying basis function. Thus, Eq. [2] can be expressed in the following matrix form in consideration of Eq. [3],

$$[K]\{p\} + [M]\{\dot{p}\} + [C]\{\ddot{p}\} = [B] \quad [4]$$

where $\{p\}_i$, $\{\dot{p}\}_i$, $\{\ddot{p}\}_i$ are the values of the pressure and its derivatives at time t and the elements of the matrix are written

$$K_{ij} = \int_s \nabla \psi_i \cdot \nabla \psi_j ds + \frac{1}{2r} \oint_l \psi_i \psi_j dl \quad [5a]$$

$$C_{ij} = \frac{1}{v_0} \oint_l \psi_i \psi_j dl \quad [5b]$$

$$M_{ij} = \int_s \psi_i \psi_j \sum_k \psi_k O_k ds \quad [5c]$$

$$B_i = \frac{\beta}{C_p} \int_s \psi_i \left(\sum_k \psi_k \Psi_k \right) ds \frac{\partial I}{\partial t} \quad [5d]$$

$$\{p\} = \{p_1, p_2, \dots, p_N\}^T \quad [5e]$$

$$\{\dot{p}\} = \{\dot{p}_1, \dot{p}_2, \dots, \dot{p}_N\}^T \quad [5f]$$

$$\{\ddot{p}\} = \{\ddot{p}_1, \ddot{p}_2, \dots, \ddot{p}_N\}^T \quad [5g]$$

Here, the Newmark's time stepping scheme is used for the discretization of the time dimension (28). To form an image from a presumably uniform initial guess of the microwave and mechanical property distributions, the iterative Newton's method is used to update Ψ and O from their starting values. In this method, we Taylor expand p about an assumed (Ψ, O) distribution, which is a perturbation away from some other distribution $(\tilde{\Psi}, \tilde{O})$, such that a discrete set of p values can be expressed as

$$p(\tilde{\Psi}, \tilde{O}) = p(\Psi, O) + \frac{\partial p}{\partial \Psi} (\tilde{\Psi} - \Psi) + \frac{\partial p}{\partial O} (\tilde{O} - O) \quad [6]$$

If we assume the microwave/mechanical property distributions are close to the true profiles, the left-hand side of Eq. [6] can be considered as the true data (observed or measured), and the relationship can be truncated to yield

$$J\Delta\chi = p^o - p^c \quad [7]$$

where

$$J = \begin{bmatrix} \frac{\partial p_1}{\partial O_1} & \dots & \frac{\partial p_1}{\partial O_K} & \frac{\partial p_1}{\partial \Psi_1} & \dots & \frac{\partial p_1}{\partial \Psi_K} \\ \frac{\partial p_2}{\partial O_1} & \dots & \frac{\partial p_2}{\partial O_K} & \frac{\partial p_2}{\partial \Psi_1} & \dots & \frac{\partial p_2}{\partial \Psi_K} \\ \vdots & \ddots & \vdots & \vdots & \ddots & \vdots \\ \frac{\partial p_M}{\partial O_1} & \dots & \frac{\partial p_M}{\partial O_K} & \frac{\partial p_M}{\partial \Psi_1} & \dots & \frac{\partial p_M}{\partial \Psi_K} \end{bmatrix} \quad [8a]$$

$$\Delta\chi = (\Delta O_1, \Delta O_2, \dots, \Delta O_K, \Delta \Psi_1, \Delta \Psi_2, \dots, \Delta \Psi_K)^T \quad [8b]$$

$$p^o = (p_1^o, p_2^o, \dots, p_M^o)^T \quad [8c]$$

$$p^c = (p_1^c, p_2^c, \dots, p_M^c)^T \quad [8d]$$

and p_i^o and p_i^c are the measured and calculated acoustic field data based on the assumed (p, Ψ) distribution data for $i=1,2,\dots,M$ boundary locations. O_k and Ψ_k ($k=1,2,\dots,K$) are the reconstruction parameters for the elastic and microwave property profiles, respectively. In order to realize an invertible system of equations for $\Delta\chi$, Eq. [7] is left multiplied by the transpose of J to produce

$$(J^T J + \lambda I) \Delta\chi = J^T (p^o - p^c) \quad [9]$$

where regularization schemes are invoked to stabilize the decomposition of $J^T J$. I is the identity matrix, and λ is the regularization scheme combined Marquardt and Tikhonov regularization schemes. We have found that when $\lambda=(p^o-p^c) \times \text{trace}[J^T J]$, the reconstruction algorithm generates best results, where $\text{trace}[\cdot]$ represents the trace of the matrix. The process now involves determining the calculated hybrid regularization-based Newton's method to update an initial microwave/elastic property distribution iteratively via the solution of Eqs. [4,9] so that an object function composed of a weighted sum of the squared difference between computed and measured acoustic pressures can be minimized. The objective function F is given by

$$\text{Min} : F = \sum_1^M (p_i^o - p_i^c)^2 \quad [10]$$

Considering the fact that US transducers and signal amplifiers have band-pass characteristics, we used the Butterworth band-pass filter module in the iterative process to improve the accuracy of reconstruction. The flowchart of the algorithm is shown in *Figure 1*. Before starting the calculation, we set microwave energy loss to 0.0001 a.u. and the elasticity coefficient to 1.90 GPa.

We conducted several simulations and phantom

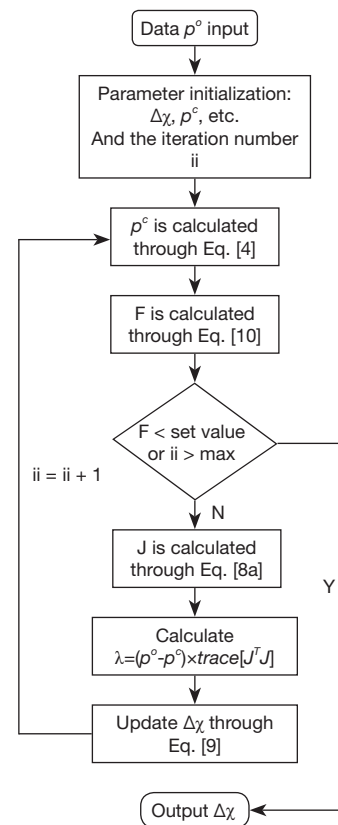


Figure 1 Flowchart.

experiments to demonstrate the feasibility of the proposed TAE.

Numerical simulations

In the first simulation experiment, different contrast levels of elastic modulus between the targets and the background and microwave energy loss were used. A circular background region with the radius of 20 mm contained three circular targets (2 mm in radius each), positioned at 9, 2 and 6 o'clock, respectively. The bulk elastic modulus and microwave energy loss of the targets are listed in *Table 1*.

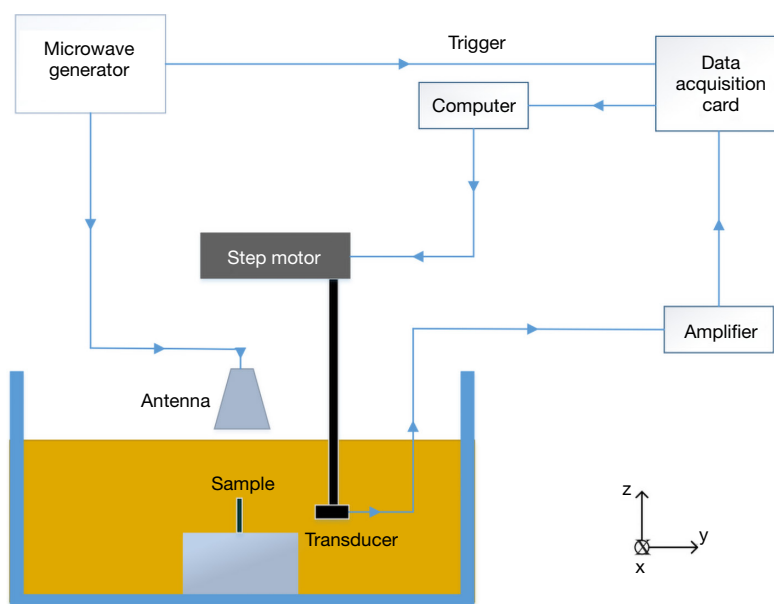
Similarly to the first simulation, a circular background region with a radius of 20 mm contained three circular targets positioned at 9, 2, and 6 o'clock with a radius of 5, 2, and 2 mm, respectively. The bulk elastic modulus, relative dielectric constant and conductivity of the targets are listed in *Table 2*. The microwave energy loss distribution was obtained by solving the electromagnetic wave equation. The microwave energy loss distribution obtained was then input into the algorithm for elastic modulus recovery. In

Table 1 Parameters of the background and targets in the first simulation

Variable	Radius (mm)	Location	Elastic modulus (GPa)	Microwave energy loss (a.u.)
Background			1.90	0.001
Target 1	2	2 o'clock	2.43	1
Target 2	2	6 o'clock	2.34	2
Target 3	2	9 o'clock	2.25	1.5

Table 2 Parameters of the background and targets in second simulation

Variable	Radius (mm)	Location	Elastic modulus (GPa)	Relative dielectric constant (a.u.)	Conductivity (S/m)
Background			1.90	2	0.001
Target 1	2	2 o'clock	2.43	56	5
Target 2	2	6 o'clock	2.34	60	5
Target 3	5	9 o'clock	2.25	66	5

**Figure 2** Schematic of the experimental setup.

the simulation setting, a plane wave at 3 GHz was incident along the x axis.

The reconstructed method was implemented based on a dual mesh scheme (30) with 1,584 nodes and 2,986 elements of coarse mesh and 6,153 nodes and 11,944 elements of fine mesh. Thirty iterations were required for a complete image reconstruction. Each iteration cost ~6 min on a 3.6 Ghz core processor with 20 GB memory.

Phantom experiments

The TAT system (19–21) we used to demonstrate the reconstruction method is schematically shown in *Figure 2*, where a coordinate system (X, Y, Z) is also depicted for reference. In this single transducer scanning experimental system, a pulsed microwave was generated from a custom-designed microwave generator (frequency: 3.0 GHz, bandwidth: 50 MHz, peak power: 70 KW, pulse duration:

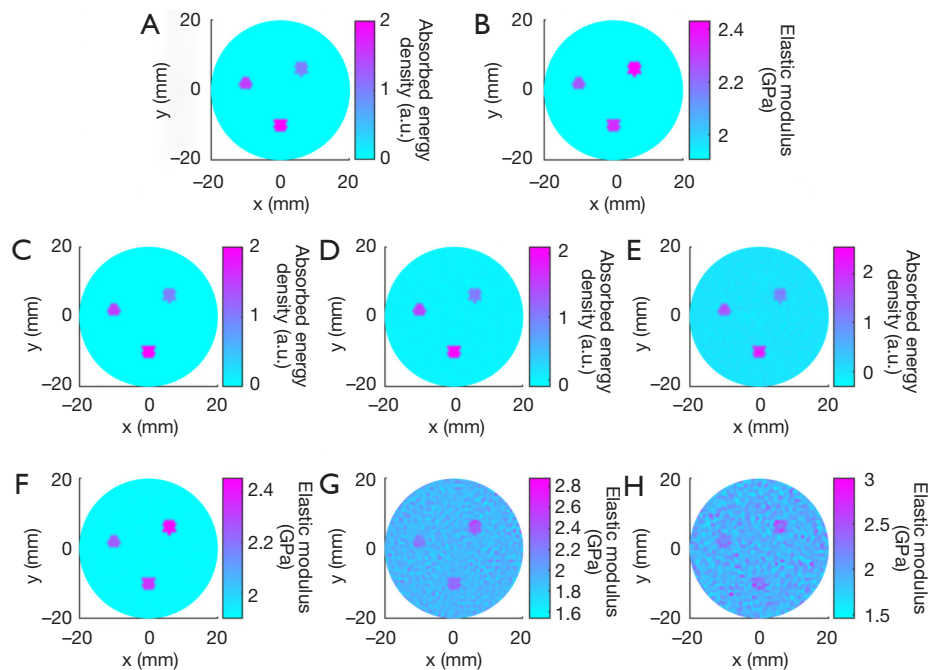


Figure 3 Reconstructed TAE images from the first simulation. Initial microwave energy loss distribution (A) and elastic modulus distribution (B). (C,D,E) The reconstructed microwave energy loss distribution for cases 1, 2, and 3, respectively. (F,G,H) The reconstructed elastic modulus distribution for cases 1, 2, and 3, respectively. TAE, thermoacoustic elastography.

750 ns). The microwave pulse was coupled into a rectangular waveguide and irradiated to a sample with less than 3.2 W/m^2 power density at the sample surface by a horn antenna ($114 \times 144 \text{ mm}^2$). A single-element unfocused transducer (V323, Olympus, central frequency: 2.25 MHz, nominal bandwidth: 1.35 MHz) was used to acquire the TA signal in a polyvinyl chloride tank filled with transformer oil for optimal microwave induced US signal transmission. TA signals at 180 positions were received via a step motor rotating the transducer around the sample over 360° in the X-Y plane. The signals were recorded by a data acquisition card (PCI4732, Vidts Dynamic) after amplified by a 58 dB amplifier (homemade, bandwidth: 250 kHz–2MHz).

In the phantom experiments, agar-based tissue-mimicking phantoms were used as the investigation samples. The elastic modulus was varied by changing the concentration of agar, while the microwave energy absorption was changed by adding 4% salt to the phantom. In the experiments, we made 2% (2.2 GPa elastic modulus), 4% (2.34 GPa elastic modulus) and 6% (2.43 GPa elastic modulus) agarose phantoms (3 mm in diameter each) with the transform oil background (1.90 GPa elastic modulus) for TAE. To improve the signal to noise (SNR), an average

of 100 times was adopted during signal acquisition. The same finite element meshes as the simulations were used for the reconstruction using the experimental data.

In the first phantom experiment, a single target having 2% or 6% agarose was placed in the background. In the second phantom experiment, two cases were examined. In case 1, there were two 4% agarose targets. In case 2, there was a 2% agarose target and a 6% agarose target.

Results

Numerical simulations

In the first numerical simulation experiment, the anti-noise ability of the algorithm was also tested. For cases 1, 2, and 3, “measured” data without noise and with 20 and 10 dB SNR were used as the reconstructed input, respectively. In each case, the measured data was generated using the forward computation with the exact elastic modulus and microwave energy loss distribution at 180 receiving/detection positions equally distributed along the boundary of the circular background region.

Figures 3 and *4* show the reconstructed images for simulation cases 1–3. We see that both the microwave

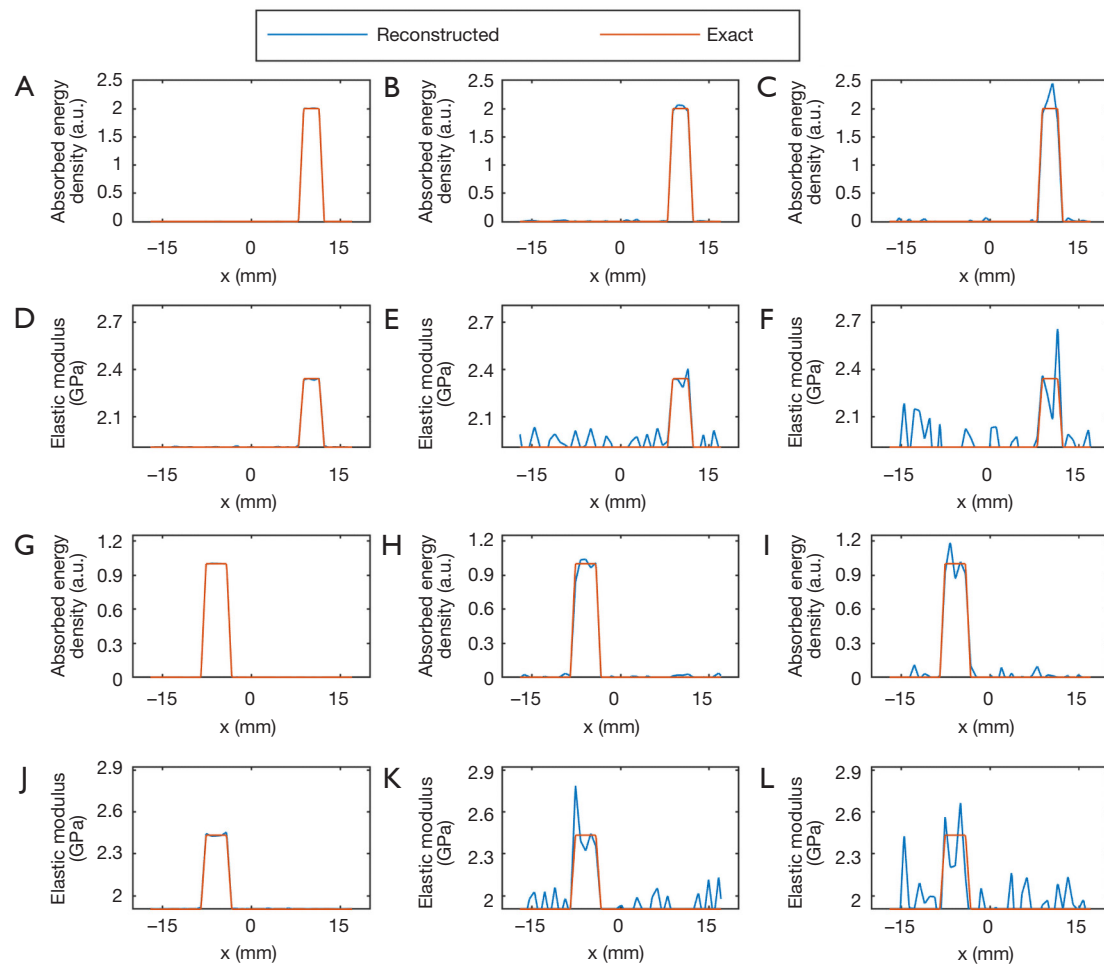


Figure 4 Comparison of the exact and reconstructed microwave energy loss (A,B,C,G,H,I) and elastic modulus (D,E,F,J,K,L) along two transects crossing the centers of targets 1 (A,B,C,D,E,F) and 2 (G,H,I,J,K,L) for the images shown in *Figure 3* for cases 1 (column 1), 2 (column 2), and 3 (column 3).

energy loss and elasticity can be well recovered for three targets having different contrast level relative to the background. When noise was added, we note that the microwave energy loss can still be reconstructed with high accuracy, while the impact of noise on the quality of elasticity image is notable. To further quantify these observations, *Figure 4* is given to present the elastic modulus and microwave energy loss profiles along the two transects crossing the centers of target 1 and target 2. We can see that recovered values of the targets match well with the exact values. Mean reconstruction error is defined as the difference between the average reconstructed value and the exact value. For microwave energy loss and elastic modulus, they were found to be 0.15%/0.46%, 3.21%/5.49%, and 9.36%/12.1% for cases 1, 2, and 3, respectively.

In the second numerical simulation, the microwave energy loss distribution was not uniform because of the difference in dielectric properties of the targets and background. For the energy loss distribution (shown in *Figure 4A*), the middle of the target at 9 o'clock collapses, and the upper part of the target at 2 and the left part of the target at 6 o'clock are missing. For this complex distribution of microwave energy loss, the anti-noise ability of the algorithm was also tested. For cases 4, 5, and 6, “measured” data without noise and with 30 and 20 dB signal-to-noise ratio (SNR) were used as the input.

Figures 5 and *6* show the reconstructed images for simulation cases 4–6. Similarly, both microwave energy loss and elastic modulus distribution are well reconstructed. We can see that recovered values of the targets match well with

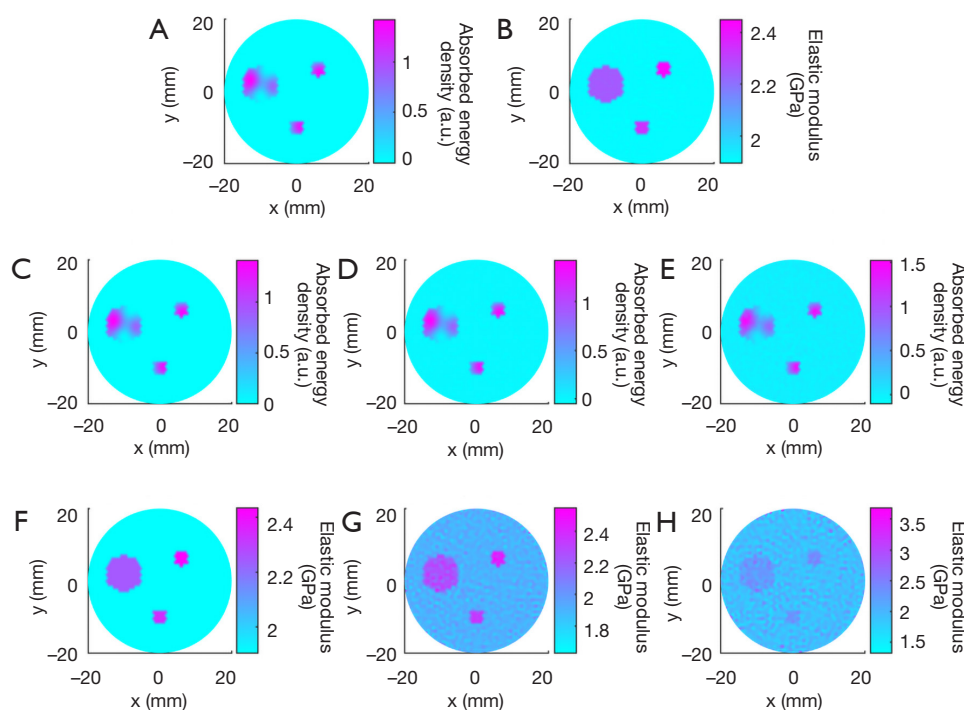


Figure 5 Reconstructed TAE images for the second simulation. Initial microwave energy loss distribution (A) and elastic modulus distribution (B). (C,D,E) The reconstructed microwave energy loss distribution for cases 4, 5, and 6, respectively. (F,G,H) The reconstructed elastic modulus distribution for cases 4, 5, and 6, respectively. TAE, thermoacoustic elastography.

the exact values. Although the energy loss distribution could not be accurately reconstructed in terms of the target shape due to the difference in dielectric constants, the recovered elastic modulus had no distortion. The mean reconstruction errors for microwave energy loss and elastic modulus were found to be 0.12%/0.11%, 2.17%/2.62%, and 5.43%/8.12% for cases 4, 5, and 6, respectively.

Phantom experiments

Figure 7 presents the reconstructed TAE images and image profiles using the phantom experiments of the first phantom experiment. In Figure 7, the elastic modulus of the first and second rows are 2.28 GPa, and the others are 2.48 GPa. From these images and image profiles, we see that our TAE algorithm has the ability to reconstruct energy density loss and elastic modulus simultaneously. The recovered shape and size of target agree well with the actual ones, and the mean error of the reconstructed elastic modulus was estimated to be less than 9.96% compared to the exact value.

The recovered TAE images (Figure 8) from the second phantom experiment are shown in Figure 8A,B,E,F, while

the reconstructed elastic modulus, microwave energy loss and elastic modulus profiles are given in Figure 8C,D,G,H. We note that the error of recovered elastic modulus for the two 2.34 GPa targets is less than 9.7% for case 1, while for case 2 the reconstruction error for the 2.43 and 2.28 GPa targets are less than 11.22% and 12.74%, respectively.

Discussion

The goal of this study is to explore a simple methodology of elastic modulus reconstruction based on conventional TAT. On the basis of the results of simulation and phantom experiments presented above, it is clear that the image reconstruction algorithm described in this paper is feasible. It has been demonstrated that the reconstructed images can provide both qualitative and quantitative information about the targets in terms of their location, sizes, shape, microwave, and mechanical property values. These observations and the quantitative nature of the reconstructed images are further confirmed by the elastic modulus profiles provided.

From Figures 3-6, we can see that both elastic modulus

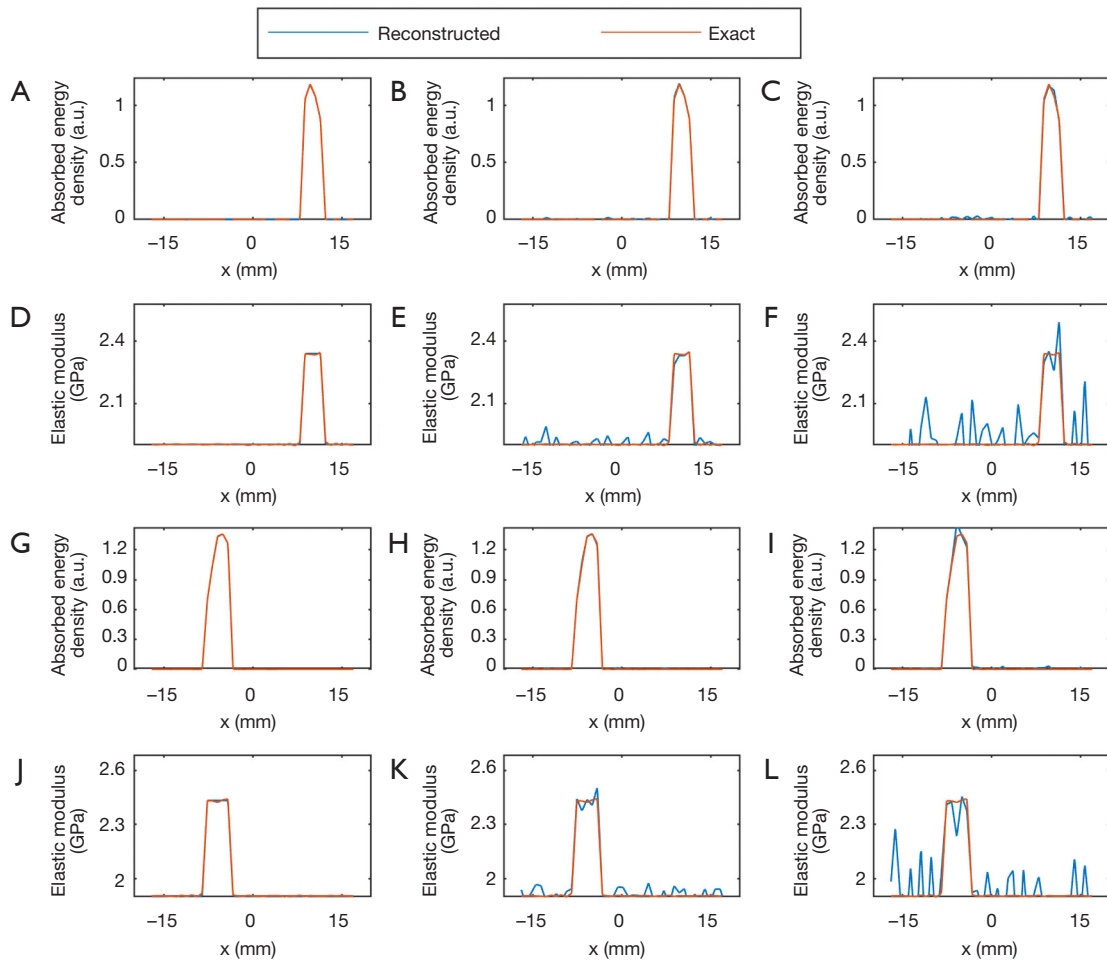


Figure 6 Comparison of the exact and reconstructed microwave energy loss (A,B,C,G,H,I) and elastic modulus (D,E,F,J,K,L) along two transects crossing the centers of targets 1 (A,B,C,D,E,F) and 2 (G,H,I,J,K,L) for the images shown in *Figure 5* for cases 4 (column 1), 5 (column 2), and 6 (column 3).

and microwave energy loss distribution of the three targets with different contrast are accurately reconstructed. We note that the accuracy for the elastic modulus reconstruction is lower than that for the microwave energy loss especially when the SNR is reduced.

The reconstructed elastic modulus given in *Figures 7* and *8* shows a good agreement with the exact value. It is worth noting that the computational time for TAE image reconstruction is long at present. We plan to speed up the reconstruction process using GPU-based parallel

computing.

Conclusions

In summary, we have demonstrated a TAE approach that allows the simultaneous recovery of elastic modulus and microwave energy loss density using both simulated and experimental data. We plan to apply this TAE approach to in vivo studies, e.g., for cirrhosis detection in the near future.

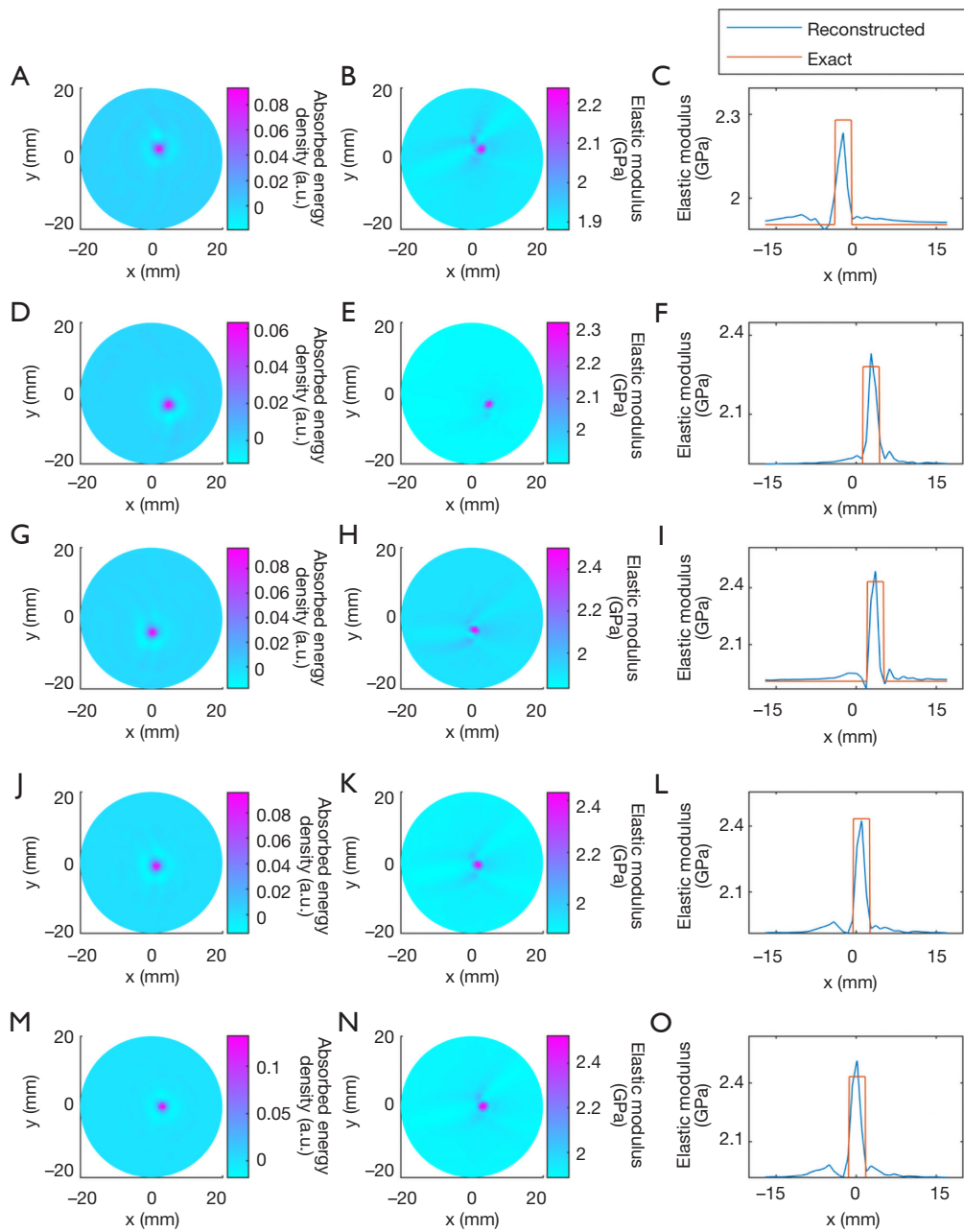


Figure 7 Reconstructed TAE images using experimental data of the first phantom experiment. Microwave energy loss density (left column), elastic modulus distribution (middle column), and the exact and reconstructed elastic modulus profiles along transects crossing the center of each target (right column). TAE, thermoacoustic elastography.

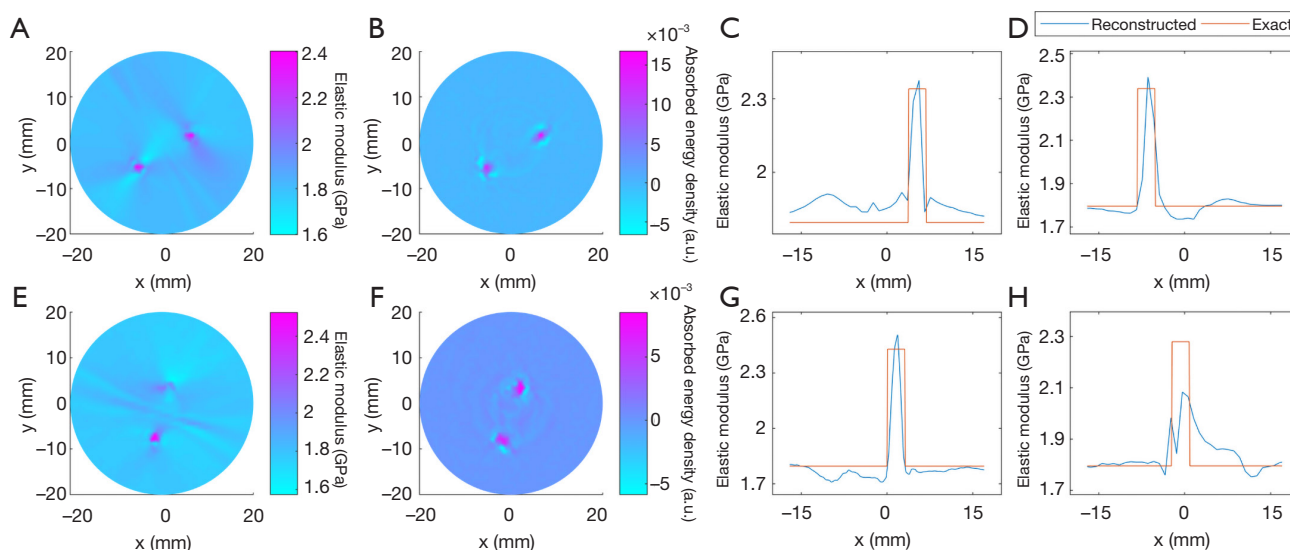


Figure 8 Reconstructed TAE images using experimental data from the second phantom experiment. The first and second rows are the images for cases 1 and 2, respectively. Elastic modulus distribution (first column), microwave energy loss density (second column), and the exact and reconstructed elastic modulus profiles along transects crossing the center of each target (the third and fourth columns). TAE, thermoacoustic elastography.

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Footnote

Conflicts of Interest: The authors have no conflicts of interest to declare.

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